Image decomposition using total variation and $\text{div}(BMO)$

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Abstract

This paper is devoted to the decomposition of an image $f$ into $u+v$, with $u$ a piecewise-smooth or “cartoon” component, and $v$ an oscillatory component (texture or noise), in a variational approach. Y. Meyer [29] proposed refinements of the total variation model (L. Rudin, S. Osher and E. Fatemi [35]) that better represent the oscillatory part $v$: the spaces of generalized functions $G = \text{div}(L^\infty)$ and $F = \text{div}(BMO)$ (this last space arises in the study of Navier-Stokes equations, H. Koch - D. Tataru [26]) have been proposed to model $v$, instead of the standard $L^2$ space, while keeping $u$ a function of bounded variation. D. Mumford and B. Gidas [27] also show that natural images can be seen as samples of scale invariant probability distributions that are supported on distributions only, and not on sets of functions. However, there is no simple solution to obtain in practice such decompositions $f = u + v$, when working with $G$ or $F$. In earlier work [45], [46], [34], the authors have proposed approximations to the $(BV, G)$ decomposition model, where the $L^\infty$ space has been substituted by $L^p$, $1 \leq p < \infty$. In the present paper, we introduce energy minimization models to compute $(BV, F)$ decompositions, and as a by-product we also introduce a simple model to realize the $(BV, G)$ decomposition. In particular, we investigate several methods for the computation of the $BMO$ norm of a function in practice. Theoretical, experimental results and comparisons to validate the proposed new methods are presented.
1 Introduction and Motivations

In what follows, we assume that a given grayscale image can be represented by a function (or sometimes distribution) \( f \), defined on an open, bounded and connected subset \( \Omega \) of \( \mathbb{R}^2 \), with Lipschitz boundary \( \partial \Omega \). In general, \( \Omega \) is a rectangle in the plane. Sometimes, we may assume that the image \( f \) is defined everywhere in the plane (obtained by extension). We limit our presentation to the two-dimensional case, but our results hold in any dimension.

We are interested in decomposing \( f \) into \( u + v \) via an energy minimization problem

\[
\inf_{(u,v) \in X_1 \times X_2} \{ K(u,v) = F_1(u) + \lambda F_2(v) : f = u + v \},
\]

where \( F_1, F_2 \geq 0 \) are functionals and \( X_1, X_2 \) are spaces of functions or distributions such that \( X_1 = \{ u : F_1(u) < \infty \}, X_2 = \{ v : F_2(v) < \infty \} \). It is assumed that \( f \in X_1 + X_2 \). The constant \( \lambda > 0 \) is a tuning parameter. Usually, \( F_1 \) and
$F_2$ are norms or semi-norms of functional spaces arising in image analysis (i.e. $F_i(\cdot) = \| \cdot \|_{X_i}$). An important problem in image analysis is to separate different features in images. For instance, in image denoising, $f$ is the observed noisy version of the true unknown image $u$, while $v$ represents additive Gaussian noise of zero mean. Often in this case, $X_1 \subset X_2$, $f \in X_2$ and $X_1$ is a space of functions “smoother” or less oscillating than those in $X_2$. However, sharp edges or boundaries have to be represented in $u$. Another related problem is the separation of the geometric (cartoon) component $u$ of $f$ from the oscillatory component $v$, representing texture or noise, of zero mean. In other cases, $u$ can be seen as a geometric or structure component of $f$, while $v$ is clutter, see [48]. A good model for $K$ is given by a choice of $X_1$ and $X_2$ so that with the above given properties of $u$ and $v$, the (semi) norms $F_1(u) = \| u \|_{X_1}$ and $F_2(v) = \| v \|_{X_2}$ are small. We give here two examples of image decomposition models by variational methods that are most related with our framework. However, many other previous work (variational or non-variational) can be seen as decompositions of $f$ into $u + v$. In the Mumford and Shah model for image segmentation [31], $f \in L^\infty(\Omega) \subset L^2(\Omega)$ is split into $u \in SBV(\Omega)$ ([30], [3], a piecewise-smooth function with its discontinuity set $J_u$ composed of a union of curves of total finite length), and $v = f - u \in L^2(\Omega)$ representing noise or texture. The problem in the weak formulation is [31], [30]

$$\inf_{(u,v) \in SBV(\Omega) \times L^2(\Omega)} \left\{ \int_{\Omega \setminus J_u} |\nabla u|^2 + \alpha \mathcal{H}^1(J_u) + \beta \|v\|^2_{L^2(\Omega)}, \ f = u + v \right\}, \quad (1)$$

where $\mathcal{H}^1$ denotes the 1-dimensional Hausdorff measure, $\alpha, \beta > 0$ are tuning parameters. With the above notations, the first two terms in the energy from (1) compose $F_1(u)$, while the third term makes $F_2(v)$. A related decomposition is obtained by the total variation minimization model of Rudin, Osher, Fatemi [35] for image denoising, where $SBV(\Omega)$ is substituted by the slightly larger space $BV(\Omega)$ of functions of bounded variation that is defined by [19], [4], [5]

**Definition 1.** Let $u \in L^1(\Omega)$; we say that $u$ is a function of bounded variation in $\Omega$ if the distributional derivative of $u$ is representable by a finite Radon measure in $\Omega$, i.e. if

$$\int_{\Omega} u \frac{\partial \phi}{\partial x_i} \, dx = - \int_{\Omega} \phi D_i u, \forall \phi \in C^1_c, \ i = 1, 2,$$

for some $\mathbb{R}^2$-valued measure $Du = (D_1 u, D_2 u)$ in $\Omega$. The vector space of all functions of bounded variation in $\Omega$ is denoted by $BV(\Omega)$.

Another equivalent definition of the space $BV(\Omega)$ (as a dual space) is obtained by:
Definition 2. Let $u \in L^1(\Omega)$. The variation of $u$ in $\Omega$ is defined by

$$V(u, \Omega) := \sup \left\{ \int_{\Omega} u(\text{div}\tilde{g})\,dx : \tilde{g} \in \left[ C^1_c(\Omega) \right]^2, \|\tilde{g}\|_{L^\infty(\Omega)} \leq 1 \right\}.$$  

Proposition 1. Let $u \in L^1(\Omega)$. Then $u \in BV(\Omega)$ if and only if $V(u, \Omega) < \infty$. In addition, $V(u, \Omega) = |Du|(\Omega)$ for any $u \in BV(\Omega)$.

Note that when $u \in W^{1,1}(\Omega)$, then $Du = \nabla u\,dx$, but the inclusion $W^{1,1}(\Omega) \subset BV(\Omega)$ is strict. However, by slight abuse of notation, we will sometimes use $|Du|(\Omega) = \int_{\Omega} |\nabla u|\,dx = |u|_{BV(\Omega)}$ for $u \in BV(\Omega)$, where $|u|_{BV(\Omega)}$ is the seminorm. Equipped with $\|u\|_{BV(\Omega)} := |Du|(\Omega) + \|u\|_{L^1(\Omega)}$, $BV(\Omega)$ becomes a Banach space.

The Rudin-Osher-Fatemi decomposition model can be defined as [35],

$$\inf_{(u,v) \in BV(\Omega) \times L^2(\Omega)} \left\{ J(u, v) = |u|_{BV(\Omega)} + \lambda \|v\|_{L^2(\Omega)}^2, \ f = u + v \right\}, \quad (2)$$

where $\lambda > 0$ is a tuning parameter. In the original TV model, $v$ represents additive Gaussian noise of zero mean. This model provides a unique $(BV(\Omega), L^2(\Omega))$ decomposition of $f \in L^2(\Omega)$, for each $\lambda > 0$ (see [10], or [44] for a more general case). The model is convex, easy to solve in practice, and denoises well piecewise-constant images while preserving edges. However it has some limitations. For instance, if $f$ is the characteristic function of a smooth set $E$ of finite perimeter, the model should produce $u = f$, $v = 0$. But this is not true for any finite value of $\lambda$, [29], [40], [5]. Cartoon or $BV$ pieces of $f$ are sent to $v$, and the model does not always represent well texture or oscillatory details, as we will see later. In [41], the authors have proposed a hierarchical multiscale $(BV(\Omega), L^2(\Omega))$ decomposition to reduce such artifacts. Also, in [24], [2], it has been shown that natural images are not well represented by functions of bounded variation.

We recall the following definition:

Definition 3. Let $\mathcal{D}(\Omega)$ be the set of test functions in $\Omega$, i.e. the set of all functions $\phi$ on $\Omega$ that are infinitely differentiable and, together with all their derivatives, are rapidly decreasing (i.e. remain bounded when multiplied by arbitrary polynomials) near the boundary $\partial\Omega$. The set of all distributions (linear continuous functionals on $\mathcal{D}(\Omega)$) is denoted by $\mathcal{D}'(\Omega)$.

Here we are interested in a better choice for the oscillatory component $v$ or for the space $X_2$, which has to give small norms for oscillatory functions, while
keeping $X_1 = BV(\Omega)$. Our discussion follows Y. Meyer [29], together with the motivations from D. Mumford and B. Gidas [27]. The idea is to use weaker norms for the oscillatory component $v$, instead of the $L^2(\Omega)$ norm, and this can be done by the use of generalized functions. For instance, Y. Meyer suggests the use of $v \in (BV(\Omega))'$, the dual of the $BV(\Omega)$ space, having the inclusions $BV(\Omega) \subset L^2(\Omega) \subset (BV(\Omega))'$. However, there is no known integral representation of continuous linear functionals on $BV(\Omega)$. There is a result that describes the dual of the $SBV(\Omega)$ space by T. De Pauw [18], but it leads to a complicated representation. To overcome this, Y. Meyer [29] suggests to approximate $(BV(\Omega))'$ by another slightly larger space, the dual $(W^{1,1}_0(\Omega))' = W^{-1,\infty}(\Omega)$. This is equivalent with the following space of distributions [1], [29]:

**Definition 4.** Let $G = G(\Omega)$ consists of distributions $T$ in $\mathcal{D}'(\Omega)$ which can be written as

$$T = div(\vec{g}) \text{ in } \mathcal{D}'(\Omega), \quad \vec{g} = (g_1, g_2) \in \left(L^\infty(\Omega)\right)^2,$$

i.e. $T(\phi) = - \int_\Omega \left(g_1 \frac{\partial \phi}{\partial x} + g_2 \frac{\partial \phi}{\partial y}\right)$ for any $\phi \in \mathcal{D}(\Omega)$.

Define $\|\cdot\|_G$ on $G$ by

$$\|T\|_G = \inf \left\{ \|\sqrt{(g_1)^2 + (g_2)^2}\|_{L^\infty(\Omega)} : T = div(\vec{g}) \text{ in } \mathcal{D}'(\Omega), \quad \vec{g} \in L^\infty(\Omega, \mathbb{R}^2) \right\}.$$

We recall that $W^{1,1}_0(\Omega)$ is the closure of $C_0^\infty(\Omega)$ in the space $W^{1,1}(\Omega)$. Functions in $W^{1,1}_0(\Omega)$ have zero trace on $\partial \Omega$. The space $G$ is a Banach space, because it is isometrically isomorphic with the dual space (equipped with the dual norm) of the normed space $W^{1,1}_0(\Omega)$ equipped with $\|u\|_{W^{1,1}_0(\Omega)} = \int_\Omega |\nabla u|$. We denote this dual space in the usual way by $W^{-1,\infty}(\Omega)$.

**Remark 1.** When $\Omega$ is open, bounded and connected in $\mathbb{R}^2$ with Lipschitz boundary, under the additional (technical) assumption $f \in L^2(\Omega)$, we have: if $f$ is decomposed into $u + v$ by a $(BV, X_2)$ model, with $u \in BV(\Omega) \subset L^2(\Omega)$ and $v = f - u \in X_2(\Omega)$, then we must have $v \in L^2(\Omega)$. Since $v$ corresponds to additive noise+texture of zero mean $\int_\Omega v = 0$, G. Aubert - J.-F. Aujol in [6], also following [29], consider the subspace $X_2 = \{v \in L^2(\Omega) : \int_\Omega v = 0\}$ of both $L^2(\Omega)$ and $G(\Omega)$ which coincides with the space $\{v = div(\vec{g}) : \vec{g} \cdot \vec{n} = 0\}$. However, the minimizers given by the $(BV(\Omega), G(\Omega))$ model will be different from the minimizers given by the ROF model [35].
We would like now to introduce the space $F$ (proposed earlier in [26] for the Navier-Stokes equations and in [29] as a suitable space for modeling textures instead of the space $G$). Let us work for a moment on the entire space $\mathbb{R}^2$ (assuming for instance that the data $f$ is extended by zero or by reflection outside the open rectangle $\Omega$).

**Definition 5.** (John-Nirenberg space of bounded mean oscillation) Let $f \in L^1_{\text{loc}}(\mathbb{R}^2)$. We say that $f$ belongs to $\text{BMO}(\mathbb{R}^2)$, if the inequality

$$\frac{1}{|Q|} \int_Q |f - f_Q| \leq A$$

holds for all squares $Q$ (it is sufficient to consider squares with sides parallel with the axis). Here $f_Q = |Q|^{-1} \int_Q f(x, y)$ denotes the mean value of $f$ over the square $Q$. The smallest such $A$ is chosen to be the norm of $f$ in $\text{BMO}(\mathbb{R}^2)$, denoted by $||f||_{\text{BMO}(\mathbb{R}^2)}$, i.e.

$$||f||_{\text{BMO}(\mathbb{R}^2)} = \sup_{Q=\text{square}} \frac{1}{|Q|} \int_Q |f - f_Q|. \quad (3)$$

Often in harmonic analysis, the Hardy space $H^1(\mathbb{R}^2)$ [37] is preferred instead of $L^1(\mathbb{R}^2)$ (with $H^1(\mathbb{R}^2) \subset L^1(\mathbb{R}^2)$), because $H^1(\mathbb{R}^2)$ has a predual which is the space $\text{VMO}(\mathbb{R}^2)$ (vanishing mean oscillation space [37]) while $L^1(\mathbb{R}^2)$ does not. When this substitution is applied, $L^\infty(\mathbb{R}^2) = \left( L^1(\mathbb{R}^2) \right)'$ is substituted by $\text{BMO}(\mathbb{R}^2) = \left( H^1(\mathbb{R}^2) \right)'$. Therefore, we are lead to consider in a similar way the space $F$ of generalized functions defined as (Y. Meyer [29], H. Koch-D. Tataru [26]):

**Definition 6.** Let $F$ consists of generalized functions $T$ which can be written as

$$T = \text{div}(\vec{g}), \quad \vec{g} = (g_1, g_2) \in \text{BMO}(\mathbb{R}^2, \mathbb{R}^2).$$

Define $||.||_F$ on $F$ by,

$$||T||_F = \inf \left\{ \left( ||g_1||_{\text{BMO}(\mathbb{R}^2)} + ||g_2||_{\text{BMO}(\mathbb{R}^2)} \right) : T = \text{div}(\vec{g}), \quad \vec{g} = (g_1, g_2) \in \text{BMO}(\mathbb{R}^2, \mathbb{R}^2) \right\}.$$
Similar with the case of the space $G$, this space $F$ can also be identified with the dual of the $H^1(\mathbb{R}^2)$-Sobolev space $L_1^1(H^1)$ (with $I_1$ the Riesz potential), see [38], or sometimes denoted by $F^1_1(\mathbb{R}^2)$ see [43]. Y. Meyer suggests that the space $F$ can also better model the oscillatory component $v$ in the $u + v$ decomposition model than the $L^2$ space, and we will show this statement in this paper. In the rest of the paper, we will work with local versions $F(\Omega)$, since $BMO(\Omega)$ is well defined also on bounded domains (even if the Hardy space $H^1$ is not well suited on bounded domains):

**Definition 7.** Let $F(\Omega)$ consists of generalized functions $T$ in $\mathcal{D}'(\Omega)$ which can be written

$$T = \text{div}(\vec{g}), \quad \vec{g} = (g_1, g_2) \in BMO(\Omega, \mathbb{R}^2),$$

equipped with the norm

$$\|T\|_{F(\Omega)} = \inf \left\{ (\|g_1\|_{BMO(\Omega)} + \|g_2\|_{BMO(\Omega)}) : T = \text{div}(\vec{g}), \quad \vec{g} = (g_1, g_2) \in BMO(\Omega, \mathbb{R}^2) \right\}.$$ 

The next two examples show why the choice of $X_2 = L^2$ or $X_2 = L^1$ does not always model oscillatory functions very well, and the proposed models, obtained by substituting $\|v\|_{L^2}$ by $\|v\|_F$ or $\|v\|_G$ in the ROF model, give better decompositions.

**Example 1.** Let $a > 0$, $n > 0$ be fixed, and let $\varphi$ be a smooth function defined on $\mathbb{R}$ such that

$$\varphi(x) = \begin{cases} a & \text{if } |x| < n, \\ 0 & \text{if } |x| > n + 1, \end{cases}$$

and $\varphi$ is increasing on $(-\infty, 0]$ and decreasing on $[0, \infty)$. Let $m > 0$, $f(x) = \frac{1}{m}\varphi'(x)\sin(mx) + \varphi(x)\cos(mx)$ and $g(x) = \frac{\varphi(x)}{m}\sin(mx) + c$, then $f = g'$.

- We have $\|f\|_G = \frac{a}{m}$. Note that $\|f\|_F < 2\|f\|_G \to 0$, as $m \to \infty$ (the inequality $\|f\|_F < 2\|f\|_G$ will be seen later).

- $\|f\|_{L^2}^2 \geq 2a^2 \int_0^n |\cos(mx)|^2 \, dx = a^2(n + \frac{1}{2m}\sin(2mn)) \to a^2n > 0$ as $m \to \infty$.

- Let $N = \lfloor \frac{mn}{2\pi} \rfloor$ be the number of complete periods of $\cos(mx)$ in the interval $[0, n]$. We may assume $N \geq 1$. Then $\|f\|_{L^1} \geq 2a \int_0^n |\cos(mx)| \, dx \geq 8aN \int_0^{\frac{n}{m}} \cos(mx) \, dx = \frac{8aN}{m} \approx 4an/\pi$. 

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Therefore, an oscillatory function has small $G$ and $F$ norms which do not depend on the domain $\Omega = (-n, n)$ and approach 0 as the frequency of oscillations increases, but with important, not so small, $L^2$ and $L^1$ norms.

**Example 2.** Let $D = D(0, R) \subset \mathbb{R}^2$ be a disk centered at the origin with radius $R$. For some $\alpha > \epsilon > 0$, consider $f = \alpha \chi_D$, $u_\epsilon = (\alpha - \epsilon) \chi_D$, and $v_\epsilon = f - u_\epsilon = \epsilon \chi_D$. If we evaluate and compare the ROF energy for two candidate solutions, $(u = f, v = 0)$ and $(u_\epsilon, v_\epsilon = f - u_\epsilon)$, we would like to have for this $f$, for any $\epsilon$,

$$
|f|_{BV} = J(f, 0) \leq |u_\epsilon|_{BV} + \lambda ||v_\epsilon||^2_{L^2},
$$

or

$$
2\pi R\alpha \leq 2\pi R\alpha + \epsilon R(\epsilon R\lambda - 2),
$$

if and only if $\lambda \geq \frac{2}{\epsilon R}, \forall \epsilon$, i.e. when $\lambda = \infty$. However, if $||.||_{L^2}^2$ is replaced by $||.||_F$, then we have for any $\epsilon$

$$
|f|_{BV} = J(f, 0) \leq |u_\epsilon|_{BV} + \lambda ||v_\epsilon||_F \leq |u_\epsilon|_{BV} + 2\lambda ||v_\epsilon||_G
$$

$$
\Leftrightarrow 2\pi \alpha R \leq 2\pi R\alpha + \epsilon R(\lambda - 2\pi),
$$

if and only if $\lambda > 2\pi$ which does not depend on $\epsilon$ and $R$ (to obtain $||v_\epsilon||_G = \epsilon \frac{R}{2}$, we use [29], Lemma 6, page 36).

This shows a limitation of the Rudin-Osher-Fatemi model: if $f = \alpha \chi_D$ (an image free of noise, piecewise-constant and with smooth discontinuity set of finite length), then the minimizer $u_\lambda$ cannot be $f$, for any finite $\lambda$. Related remarks have been made in [29], [39], [40], [5]. Moreover, suppose we decompose $f$ by the energy minimization

$$
\inf_{(u, v)} \left\{ \mathcal{K}(u, v) = |u|_{BV} + \lambda ||v||^p_{X_2} : f = u + v, p > 0 \right\},
$$

where $\| \cdot \|_{X_2}$ is a norm or a quasi-norm, and such that $\|\chi_D\|_{X_2} \neq 0$. If we start with $f = \alpha \chi_D$, then the recovered image should be $f$ for some finite $\lambda > C$. We have the following: $\mathcal{K}(f, 0) \leq \mathcal{K}(u_\epsilon, v_\epsilon)$ if and only if $p \leq 1$. Therefore, for any $p > 1$, we cannot obtain $u = f$, for any finite value of $\lambda$. We refer the reader to Cheon, Paranjpye, Vese, Osher [15], and to Chan, Esedoglu [11], and Chan, Esedoglu, Nikolova [12] for the $(BV(\Omega), L^1(\Omega))$ version of the ROF model in two dimensions in the continuous setting, which are also improvements over the original ROF model (in [15], the authors have also considered the case when in the ROF model, the fidelity term $\|f - u\|_{L^2(\Omega)}^2$ has been replaced by $\|f - u\|_{L^1(\Omega)}$; this choice gives very good reconstruction results in practice).
From such motivations, Y. Meyer [29] proposed a decomposition of \( f \), with \( X_1 = BV(\Omega) \), via
\[
\inf \left\{ \mathcal{K}(u, v) = |u|_{BV(\Omega)} + \lambda ||v||_{X_2} \right\},
\]
where the infimum is taken over \( u \in BV(\Omega) \) and \( v \in X_2 \), such that \( f = u + v \).

Here \((X_2, \| \cdot \|_{X_2})\) is either \((G(\Omega), \| \cdot \|_{G(\Omega)})\) or \((F(\Omega), \| \cdot \|_{F(\Omega)})\). However, these minimization models cannot be directly solved in practice: there is no standard calculation of the associated Euler-Lagrange equation, as it is for the ROF model which can be solved easily by finite differences.

In [45], [46], the second author together with S. Osher proposed a method to overcome the difficulty of computing \( \| \cdot \|_G \). This has been done by the energy minimization problem
\[
\inf_{u, g_1, g_2} \left\{ G_p(u, g_1, g_2) = |u|_{BV(\Omega)} + \mu \|f - u - \partial_x g_1 - \partial_y g_2\|^2_{L^p(\Omega)} + \lambda \sqrt{g_1^2 + g_2^2} \right\}.
\]

By this model, \( f \) is decomposed into \( u + v + w \), and as \( \mu \to \infty \) and \( p \to \infty \), the model approaches Meyer’s \((BV, G)\) model. The space \( G = W^{-1,\infty}(\Omega) \) is approximated by \( W^{-1,p}(\Omega) \), with \( p < \infty \) (when \( p = 2 \), \( v \) belongs to the dual of the Sobolev space \( H_0^1(\Omega) \)).

In [34], the second author together with S. Osher and A. Sole, proposed a simplified approximated method corresponding to the case \( p = 2 \). Let \( \tilde{\gamma} = \nabla P + \tilde{Q} \) be the unique Hodge decomposition of \( \tilde{\gamma} \in L^\infty(\Omega, \mathbb{R}^2) \). Using \( f - u = v = \text{div}(\tilde{\gamma}) = \triangle P \), i.e. \( P = \triangle^{-1}(f - u) \), Meyer’s \((BV, G)\) model is then approximated by
\[
\inf_u \left\{ G_2(u) = \int_\Omega |\nabla u| + \lambda \int_\Omega |\nabla(\triangle^{-1}(f - u))|^2 \right\}.
\]
This model gives an exact decomposition \( f = u + v \), with \( u \in BV(\Omega) \) and \( v \in W^{-1,2}(\Omega) = \left( H_0^1(\Omega) \right)' \), and the minimization problem has been solved using a fourth-order non-linear PDE.

In the present paper, we propose a new method to approximate Y. Meyer’s \((BV, F)\) model. We also introduce an equivalent definition of the \( BMO \) norm, using an open set formulation, which is easily formulated and computed using curve evolution technique. As a by-product, we also propose a new method for solving the \((BV, G)\) model, different from the one proposed in [7], [6].
As we have mentioned, D. Mumford and B. Gidas [27] show that natural images, as samples from scale-invariant probability distributions, cannot be modeled by functions, but instead by generalized functions, i.e. distributions in $D'(\Omega)$.

Other related models for image decomposition into cartoon and texture have been proposed recently. We mention Daubechies and Teschke [17], Starck, Elad, and Donoho [36] for variational and wavelets approaches.

In particular, we refer to Aujol et al [7], and Aubert-Aujol [6] for more properties of the space $G$ both in theory and practice, and to another approximation of the Meyer’s $(BV, G)$ model on bounded domains. We also refer to Aujol and Chambolle [8] for properties of norms that are dual to negative Sobolev and Besov norms. Our theoretical framework extends some of the results presented in Aubert-Aujol [6] for the space $G(\Omega)$ to the case of the space $F(\Omega)$.

Other related works are by Esedoglu, Osher [20], Osher-Scherzer [32], Obereder, Osher and Scherzer [33], Goldfarb, Yin [23], and among others.

We believe that the case $(BV, F)$ has not been considered in theory or in practice previously in image analysis, therefore our contribution is new also from this point of view.

The theoretical work of Koch and Tataru [26] (mentioned by Y. Meyer in [29]) uses the space $\text{div}(\text{BMO}(\mathbb{R}^2))$ for solutions of the Navier-Stokes equations. Finally, in [9], Bourgain-Brezis analyze the equation $f = \text{div}(\vec{y})$ in some limiting cases, and applications of such results can be found in Aubert-Aujol [6] for the analysis of the space $G(\Omega)$.

## 2 Definitions and Properties of the BMO Space

Here we would like to review the definitions and some basic properties of the space $\text{BMO}$. We refer the readers to “Harmonic Analysis” by E.M. Stein [37], and also to [42], [25], and [21].

Let $\Omega$ be an open and bounded subset of $\mathbb{R}^n$. For planar images, we may assume that $\Omega = (0,1) \times (0,1) \subset \mathbb{R}^2$. To simplify the notations, we will often write $BV$, $G$, $F$, ..., instead of $BV(\Omega)$, $G(\Omega)$, $F(\Omega)$, etc.

**Definition 8.** Let $f \in L^1_{\text{loc}}$. We say that $f$ belongs to $\text{BMO}^\beta$, if the inequality

$$\frac{1}{|O|} \int_O |f - f_O| \, dx \leq A_1$$

holds for the family $\mathcal{F}^\beta$ of open sets $O \subset \Omega$ such that there exist cubes $Q_1$, and $Q_2$ with $Q_1 \subset O \subset Q_2 \subset \Omega$, and $\frac{|Q_2|}{|Q_1|} \leq \beta$; here $1 \leq \beta < \infty$ is a constant,
and $f_O = \frac{|O|^{-1}}{O} \int_O f \, dx$. The smallest such $A_1$ is chosen to be the norm of $f$ in $BMO^\beta$, denoted by $||f||_{BMO^\beta}$, i.e.

$$||f||_{BMO^\beta} = \sup_{O \in \mathcal{F}^\beta} \frac{1}{|O|} \int_O |f - f_O| \, dx. \quad (5)$$

**Definition 9.** Let $f \in L^1_{\text{loc}}$. We say that $f$ belongs to $BMO$, if the inequality

$$\frac{1}{|Q|} \int_Q |f - f_Q| \, dx \leq A_2$$

holds for all cubes $Q \subset \Omega$ with sides parallel with the axes. The smallest such $A_2$ is chosen to be the norm of $f$ in $BMO$, denoted by $||f||_{BMO}$, i.e.

$$||f||_{BMO} = \sup_{Q \subset \Omega} \frac{1}{|Q|} \int_Q |f - f_Q| \, dx. \quad (6)$$

$||.||_{BMO}$ is a semi-norm vanishing on constant functions. If we identify functions in $BMO$ which are different almost everywhere by a constant, then $BMO$ becomes a Banach space. We obtain an equivalent norm if the family of cubes is replaced by the family of balls. Moreover, as mentioned in [37],

$$||f||_{BMO_p} = \left[ \sup_{Q \subset \Omega} \frac{1}{|Q|} \int_Q |f - f_Q|^p \, dx \right]^{\frac{1}{p}}, \quad (7)$$

gives an equivalent $BMO$ norm for $p \geq 1$. Here we will consider the cases $p = 1$ and $p = 2$.

**Definition 10.** A dyadic cube is a cube of the special form,

$$Q = \left\{ k_j2^{-m} < x_j < (k_j + 1)2^{-m}; 1 \leq j \leq n \right\}, \quad (8)$$

where $m$ and $k_j$, $1 \leq j \leq n$, are integers. We say $f$ has bounded dyadic mean oscillation, $f \in BMO_d$, if

$$||f||_{BMO_d} = \sup_{Q \subset \Omega \text{dyadic}} \frac{1}{|Q|} \int_Q |f - f_Q| \, dx < \infty.$$

Let $T_\alpha f(x) = f(x - \alpha)$. We say $f \in BMO_{d,\alpha}$ if $T_\alpha f \in BMO_d$.

Let $A = \{\alpha = (\alpha_1, ..., \alpha_n) : \alpha_i \in \{0, \frac{1}{2}\}\}$. Note $|A| = 2^n$. 

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Lemma 1 ([21] and [28]). $f \in BMO$ if and only if $f \in BMO_{d,\alpha}$ for all $\alpha \in A$. In fact, $||f||_{BMO} \leq 12\max_{\alpha \in A} \{||f||_{BMO_{d,\alpha}}\}$.

Remark 2. Let $BMO_D = \bigcap_{\alpha \in A} BMO_{d,\alpha}$ with

$$||\cdot||_{BMO_D} = \max_{\alpha \in A} \{||\cdot||_{BMO_{d,\alpha}}\}.$$ 

Then the above lemma shows that $BMO = BMO_D$ with equivalent norms.

Lemma 2. Let $f \in L^1_{loc}$ and $c \in \mathbb{R}$, then

$$\int_O |f - f_O| \, dx \leq 2 \int_O |f - c| \, dx,$$

for all $O \subset \Omega$.

Proof:

$$|f - f_O| \leq |f - c| + |c - f_O| = |f - c| + |c - |O||^{-1} \int_O f \, dx$$

$$\leq |f - c| + |O|^{-1} \int |c - f| \, dx.$$

Integrating both sides over $O$, we obtain (9). Moreover,

$$\inf_{c \in \mathbb{R}} \int_O |f - c| \, dx \leq \int_O |f - f_O| \, dx \leq 2 \inf_{c \in \mathbb{R}} \int_O |f - c| \, dx.$$

Lemma 3. $||.||_{BMO^\beta}$ and $||.||_{BMO}$ are equivalent, for any $1 \leq \beta < \infty$.

Proof:

It suffices to show that there exist constants $c_1$ and $c_2$ greater than 0 such that, for all $f \in BMO$,

$$c_1 ||f||_{BMO} \leq ||f||_{BMO^\beta} \leq c_2 ||f||_{BMO}.$$

(10)

It is clear that the first inequality in (10) holds with $c_1 = 1$. It remains to show $||f||_{BMO^\beta} \leq c_2 ||f||_{BMO}$, for some $c_2 > 0$. 

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Let \( O \in \mathcal{F}_\beta \). There exist \( Q_1 \) and \( Q_2 \) in \( \Omega \) such that \( Q_1 \subset O \subset Q_2 \), and \( \frac{|Q_2|}{|Q_1|} \leq \beta \). We have

\[
\frac{1}{|O|} \int_O |f - f_O| \, dx \leq 2 \inf_{c \in \mathbb{R}} \frac{1}{|O|} \int_O |f - c| \, dx \leq 2 \inf_{c \in \mathbb{R}} \frac{1}{|Q_1|} \int_{Q_2} |f - c| \, dx
\]

\[
\leq 2 \beta \inf_{c \in \mathbb{R}} \frac{1}{|Q_2|} \int_{Q_2} |f - c| \, dx \leq 2 \beta \left( \frac{1}{|Q_2|} \int_{Q_2} |f - f_{Q_2}| \, dx \right)
\]

\[
\leq 2 \beta \sup_{Q \subset \Omega} \frac{1}{|Q|} \int_Q |f - f_Q| \, dx = ||f||_{BMO}.
\]

Taking the supremum over all \( O \in \mathcal{F}_\beta \), we have \( ||f||_{BMO} \leq c_2 ||f||_{BMO} \) with \( c_2 = 2 \beta \).

The next simple property shows that functions in \( BMO \) are scale invariant. For simplicity, assume here that \( \Omega = \mathbb{R}^n \).

**Lemma 4.** \( f(x) \) and \( f(\alpha x) \) have the same norm in \( BMO \), for all \( \alpha > 0 \).

**Proof:** Using a change of variable by letting \( y = \alpha x \), we have

\[
\frac{1}{|Q|} \int_Q \left| f(\alpha x) - \frac{\int_Q f(\alpha x) \, dx}{|Q|} \right| \, dx = \frac{1}{|Q|} \int_{\alpha Q} \left| f(y) - \frac{\int_{\alpha Q} f(y) \, dy}{|\alpha Q|} \right| \frac{1}{\alpha^n} \, dy,
\]

\[
= \frac{1}{|\alpha Q|} \int_{\alpha Q} \left| f(y) - \frac{\int_{\alpha Q} f(y) \, dy}{|\alpha Q|} \right| \, dy,
\]

\[
= \frac{1}{|Q'|} \int_{Q'} \left| f(y) - \frac{\int_{Q'} f(y) \, dy}{|Q'|} \right| \, dy,
\]

(11)

where \( Q' = \alpha Q = \{ \alpha x : x \in Q \} \); it is clear that

\[
\sup_{Q \subset \mathbb{R}^n} \frac{1}{|Q|} \int_Q \left| f(x) - \frac{\int_Q f(x) \, dx}{|Q|} \right| \, dx = \sup_{Q' \subset \mathbb{R}^n} \frac{1}{|Q'|} \int_{Q'} \left| f(y) - \frac{\int_{Q'} f(y) \, dy}{|Q'|} \right| \, dy.
\]

Note that the norms on \( G \) and \( F \) are not scale invariant, but satisfy the scaling relation \( \|f(\alpha \cdot)\|_{F,G} = \frac{1}{\alpha} \|f(\cdot)\|_{F,G} \).

**Remark 3.** \( L^\infty \subset BMO^\beta = BMO \), for \( 1 \leq \beta < \infty \), with \( ||f||_{BMO} \leq 2 ||f||_{\infty} \). Note that if in the definition of \( BMO^\beta \) we do not impose any bound on \( \beta \), i.e. we
allow $\beta = \infty$, then $BMO^\beta$ approaches the space $L^\infty$: the norm is attained at a union of very small regions, such that $f$ has largest or smallest values inside these regions. $BMO$ also contains unbounded functions; indeed, $\ln(|P|) \in BMO$ for any polynomial $P$.

Some additional properties are as follows. $|f|$ is in $BMO$ whenever $f$ is, since $||f| - |f_Q|| \leq |f - f_Q|$. However, if $|f|$ is in $BMO$, then $f$ is not necessarily in $BMO$. For example in $\mathbb{R}$, $f(x) = \text{sign}(x) \ln(|x|)$ is not in $BMO(\mathbb{R})$, while $|f(x)| = \ln(|x|)$ is. Indeed for $\alpha > 0$, $f(\alpha x) = \text{sign}(x) \ln(\alpha) + \text{sign}(x) \ln(|x|)$, and if $I = [-1, 1]$, then

$$
\frac{1}{|\alpha I|} \int_{\alpha I} \left| f(y) - \frac{\int_{\alpha I} f(y) dy}{|\alpha I|} \right| dy \to \infty, \text{ as } \alpha \to \infty.
$$

3 Numerical Computation of the $BMO$ Norm

In this section, we introduce and discuss several new methods for computing or approximating the $BMO$ norm of a given function $f$ in two dimensions, using the equivalent definitions introduced in the previous section.

3.1 Computing the $BMO$ Norm Using the Open Set Formula

Let $\phi$ be a Lipschitz-continuous function that defines $\partial O$ implicitly, i.e. $O = \{x \in \Omega : \phi(x) > 0\}$, and denote by $H(\phi)$ the Heaviside function,

$$
H(r) = \begin{cases} 
1 & \text{if } r \geq 0, \\
0 & \text{otherwise}.
\end{cases}
$$

Using the variational level set formulation, as in [47], [13], [14], we define,

$$
\mathcal{G}(\phi) = \frac{1}{\int_{\Omega} H(\phi) \, dx} \int_{\Omega} \left| g - \frac{\int_{\Omega} g H(\phi) \, dx}{\int_{\Omega} H(\phi) \, dx} \right| H(\phi) \, dx.
$$

Similarly, we also have an equivalent $||\cdot||_{BMO^p}$ norm using $p = 2$:

$$
\mathcal{F}(\phi) = \left[ \frac{1}{\int_{\Omega} H(\phi) \, dx} \int_{\Omega} \left| g - \frac{\int_{\Omega} g H(\phi) \, dx}{\int_{\Omega} H(\phi) \, dx} \right|^2 H(\phi) \, dx \right]^\frac{1}{2}.
$$
Remark 4. We will incorporate in (12) the perimeter of the unknown open set $O$ as given by $\int_\Omega |\nabla H(\phi)|$, in order to insure that $O$ remains “bulky”, i.e. we do not allow $O$ to break into very small pieces. We do not exactly impose in this way the required constraint from Definition 8, but it is a way of keeping the ratio between $|O|$ and $\partial O$ bounded. This was kindly suggested to us by Jean-Michel Morel.

If $\phi$ solves $\sup_\phi \left\{ G_{\text{new}}(\phi) = G(\phi) - \lambda \int_\Omega |\nabla H(\phi)| \, dx \right\}$ for some parameter $\lambda > 0$, then $||g||_{BMO^\beta}$ is well approximated by $G(\phi)$.

Let $H_\epsilon(\phi)$ be a smoother function approximating $H(\phi)$ as $\epsilon \to 0$.

Using the notation $|O| = \int_\Omega H_\epsilon(\phi) \, dx$, we obtain,

$$\frac{\partial G_{\text{new}}}{\partial \phi} = \left[ -\frac{1}{|O|^2} (g - g_0) \int_\Omega \frac{g - g_0}{|g - g_0|} H_\epsilon(\phi) \, dx + \frac{1}{|O|} |g - g_0| - \frac{G(\phi)}{|O|} \right. + \left. \lambda \text{div} \left( \frac{\nabla \phi}{|\nabla \phi|} \right) \right] \delta_\epsilon(\phi),$$

where $\delta_\epsilon = H_\epsilon'$. By introducing an artificial time, we then solve

$$\phi_t = \frac{\partial G_{\text{new}}}{\partial \phi} \quad \text{in } \Omega, \quad \text{and } \frac{\nabla u}{|\nabla u|} \cdot \vec{n} \text{ on } \partial \Omega.$$

However, in practice, since our approximation $\delta_\epsilon(\phi) > 0$ for any $\phi$ as in [13], we neglect this factor $\delta_\epsilon(\phi)$, and we solve the equation (to obtain faster results),

$$\phi_t = -\frac{1}{|O|^2} (g - g_0) \int_\Omega \frac{g - g_0}{|g - g_0|} H(\phi) \, dx + \frac{1}{|O|} |g - g_0| - \frac{G(\phi)}{|O|} \right. + \left. \lambda \text{div} \left( \frac{\nabla \phi}{|\nabla \phi|} \right).$$

In our numerical calculations, we always have that $Q_2(O) = \Omega$ (therefore $|Q_2|$ is bounded independent of $O$), while the existence of $Q_1(O)$, with size that does not become too small, is insured by the additional length term.

We do not guarantee that we compute a global maximum of the energy. However, the numerical experiments (using piecewise constant images) show that we obtain good and stable approximations to the exact solution, as illustrated in the section of experimental results.
3.2 Computing the $BMO$ Norm Using the Square Formulation

Note that, in two dimensions, the set $Q' = \{ |x| + |y| < r \}$ is a square centered at the origin, with side length $l(Q') = \sqrt{2}r$. The corners of $Q'$ are at the vertices $(r,0), (0,r), (-r,0), (0,-r)$. Therefore, to have the sides of $Q'$ of length $r$ and parallel to the axis, we need to rotate $\frac{1}{\sqrt{2}} Q'$ by an angle of $\frac{\pi}{4}$, i.e. by applying the matrix

$$
\begin{bmatrix}
1 & -1 \\
1 & 1
\end{bmatrix}
$$

to $Q'$. In other words, $Q = \{ |(x - x_0) - (y - y_0)| + |(x - x_0) + (y - y_0)| < r \}$ is the square centered at $(x_0, y_0)$, with sides parallel to the axis and $l(Q) = r$.

Let $\phi(x,y) = r - (|(x - x_0) - (y - y_0)| + |(x - x_0) + (y - y_0)|)$ be the signed distance function to $\partial Q$, and $H(\phi)$ be the Heaviside function. Define

$$
G(Q) = G(r, x_0, y_0) = \frac{1}{\int_\Omega H(\phi)} \int_\Omega \left| g - \frac{\int_\Omega g H(\phi)}{\int_\Omega H(\phi)} \right| H(\phi). \quad (14)
$$

Then,

$$
||g||_{BMO} = \sup_{(r,x_0,y_0)} G(r,x_0,y_0).
$$

We also have an equivalent $||.||_{BMO}$ norm with $p = 2$:

$$
F(Q) = F(r, x_0, y_0) = \left[ \frac{1}{\int_\Omega H(\phi)} \int_\Omega \left| g - \frac{\int_\Omega g H(\phi)}{\int_\Omega H(\phi)} \right|^2 H(\phi) \right]^\frac{1}{2}. \quad (15)
$$

Let $H_\epsilon$ be a smooth approximation of $H$, and $\delta_\epsilon = H'_\epsilon$. We have

$$
\frac{\partial G}{\partial r} = \frac{1}{\int_\Omega H_\epsilon(\phi)} \left[ -g \int_\Omega \frac{\partial H_\epsilon(\phi)}{\partial r} + \int_\Omega |g - g_Q| \frac{\partial H_\epsilon(\phi)}{\partial r} 
- \frac{1}{\int_\Omega H_\epsilon(\phi)} \int_\Omega \frac{g - g_Q}{|g - g_Q|} H_\epsilon(\phi) \int_\Omega (g - g_Q) \frac{\partial H_\epsilon(\phi)}{\partial r} \right],
$$

$$
\frac{\partial G}{\partial x_0} = \frac{1}{\int_\Omega H_\epsilon(\phi)} \left[ -g \int_\Omega \frac{\partial H_\epsilon(\phi)}{\partial x_0} + \int_\Omega |g - g_Q| \frac{\partial H_\epsilon(\phi)}{\partial x_0} 
- \frac{1}{\int_\Omega H_\epsilon(\phi)} \int_\Omega \frac{g - g_Q}{|g - g_Q|} H_\epsilon(\phi) \int_\Omega (g - g_Q) \frac{\partial H_\epsilon(\phi)}{\partial x_0} \right],
$$

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where \( \frac{\partial H(\phi)}{\partial \epsilon} = \delta(\phi) \), and \( \frac{\partial H(\phi)}{\partial x_0} = \left[ \frac{(x-x_0)-(y-y_0)}{[(x-x_0)-(y-y_0)]} + \frac{(x-x_0)+(y-y_0)}{[(x-x_0)+(y-y_0)]} \right] \delta(\phi) \). Similarly,

\[
\frac{\partial G}{\partial y_0} = \frac{1}{\int_{\Omega} H(\phi)} \left[ - G \int_{\Omega} \frac{\partial H(\phi)}{\partial y_0} + \int_{\Omega} |g-g_Q| \frac{\partial H(\phi)}{\partial y_0} \right] - \frac{1}{\int_{\Omega} H(\phi)} \int_{\Omega} \frac{g-g_Q}{|g-g_Q|} H(\phi) \int_{\Omega} (g-g_Q) \frac{\partial H(\phi)}{\partial y_0} \right],
\]

where \( \frac{\partial H(\phi)}{\partial y_0} = \left[ \frac{(x-x_0)-(y-y_0)}{[(x-x_0)-(y-y_0)]} + \frac{(x-x_0)+(y-y_0)}{[(x-x_0)+(y-y_0)]} \right] \delta(\phi) \).

By introducing an artificial time, we will solve the equations,

\[
\frac{\partial r}{\partial t} = \frac{\partial G}{\partial r}, \quad \frac{\partial x_0}{\partial t} = \frac{\partial G}{\partial x_0}, \quad \frac{\partial y_0}{\partial t} = \frac{\partial G}{\partial y_0}.
\]

 Again, we do not show that this method converges to a global maximum of the energy. However, in the experimental results (using piecewise constant images), we have obtained the correct answer, when we know the exact solution.

If we would work with disks instead of squares, then we could have

\[
\phi(x, y) = r^2 - (x - x_0)^2 - (y - y_0)^2,
\]

which is differentiable everywhere.

### 3.3 Exact Computation of the \textit{BMO} Norm

As kindly suggested by Jean-Michel Morel, we have also implemented an exact evaluation of the \textit{BMO} norm using the square formulation. This is computationally more expensive but still it can be made relatively fast by using FFT. In addition it produces very accurate results. The procedure is as follows.

1. Fix a list of growing scales \( \sigma = 2, 4, 8, 16, \ldots \)

2. For each \( \sigma \), consider the function \( k_{\sigma}(x, y) = \frac{1}{\sigma^2} \chi_{Q_{\sigma}(x, y)} \), where \( Q_{\sigma}(x, y) \) is the square centered at \( (x, y) \) having length \( \sigma \).

3. Convolve \( f \) with \( k_{\sigma}(0,0) \), called \( f_{\sigma}(x, y) \).

4. For each \( (x, y) \), compute \( \text{osc}(x, y) = \int |f(s, t) - f_{\sigma}(x, y)| k_{\sigma}(x, y) \, dsdt \).

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5. Take the sup of \( \text{osc}(x, y) \) which yields a value of \( \| f \|_{BMO} \) at scale \( \sigma \), denoted \( \| f \|_{BMO, \sigma} \).

6. Compute the maximal value of \( \| f \|_{BMO, \sigma} \) for all \( \sigma \)'s.

If one takes \( \sigma = 2, 3, 4, 5, \ldots, N \), where \( N \) is the size of the image \( f \), one is ensured to get the exact value of \( \| f \|_{BMO} \) in a discrete framework. In this way, one obtains the global maximum and the square(s) where it is attained. This allows to compare with the maximization procedure and see whether or not it yields the same value.

Finally, we have used the above procedure to evaluate the necessary expression over the dyadic squares and the additional translations. This gives an accurate method too, and it is faster since we have fewer squares to consider.

\section{A \( (BV, F) \) Image Decomposition Model}

Recall Meyer’s model, which decomposes \( f \) into \( u + v \), by the variational problem

\[
\inf \{ \mathcal{E}(u, v) = |u|_{BV} + \lambda \| v \|_{F} \},
\]

where the infimum is taken over \( u \in BV \) and \( v \in F \), such that \( f = u + v \).

The space \( F \) is defined as

\[
F = \{ v = \text{div}(\tilde{g}) \text{ in } \mathcal{D}' : \tilde{g} \in BMO(\Omega, \mathbb{R}^2) \}.\]

\textbf{Remark 5.} Given \( f \in L^2 \), there exists \( u \in BV \), and \( v \in L^2 \subset G \subset F \), such that \( \mathcal{E}(u, v) < \infty \), \( f = u + v \), and \( \int f = \int u \).

Indeed, recall the R-O-F model, which minimizes the functional,

\[
\mathcal{J}(u, v) = |u|_{BV} + \lambda' \| v \|_{L^2}
\]

over the set of \( u \in BV \), \( v \in L^2 \), such that \( f = u + v \). Pick \( \lambda' \) so that \( \| f \|_G > \frac{1}{2\lambda'}. \) The existence of a minimizer, denoted \( (u, v) \), for the R-O-F model has been proved in [10] and characterized in [29] as being, \( u \in BV \), \( v \in G \subset F \), with \( \| v \|_G = \frac{1}{2\lambda'}. \) Therefore \( \mathcal{E}(u, v) < \infty \).

The next theorem shows the existence of minimizers for Meyer’s \( (BV, F) \) model. We refer the reader to Aubert-Aujol [6] for a similar proof when \( \| \cdot \|_* = \| \cdot \|_G \).
Theorem 1. Let $f \in L^2$. The following minimization problem

$$\inf_{(u,v)} \left\{ \mathcal{E}(u, v) = |u|_{BV} + \lambda \|v\|_F, \; \int_\Omega u = \int_\Omega f, \; f = u + v \right\}$$  \hspace{1cm} (18)$$

has at least one solution $u \in BV$, $v = f - u \in F \cap L^2$.

**Proof:** We use the standard tool in calculus of variations. Let $\{(u_n, v_n)\}$ be a minimizing sequence (from the previous remark, we know that the infimum of the energy is finite). Then $f = u_n + v_n$ and $\int_\Omega u_n = \int_\Omega f$, $\forall n \geq 0$. In addition, there is a constant $C$ (that may change from line to line) such that $|u_n|_{BV} \leq C$, $\|v_n\|_F \leq C$, uniformly.

By Poincare-Wirtinger inequality,

$$||u_n - \int_\Omega u_n ||_{L^2} \leq C |u|_{BV},$$

and since $\int_\Omega u_n = \int_\Omega f$, $\forall n$,

$$\Rightarrow \|u_n\|_{L^2} \leq C,
\Rightarrow |u_n|_{BV} \leq C.$$ 

Then $\|u_n\|_{L^1} \leq C$ since $\Omega$ is bounded; furthermore, $\|u_n\|_{BV} = \|u_n\|_{L^1} + |u_n|_{BV} \leq C$. Therefore, there exists $u \in BV$ and a subsequence (still denoted by $u_n$), such that $u_n$ converges to $u$ in the $BV$-weak* topology. In particular, $u_n$ converges to $u$ strongly in $L^1$ and by the lower semicontinuity of the total variation, $|u|_{BV} \leq \liminf_{n \to \infty} |u_n|_{BV}$.

As for the subsequence $v_n$, we have $v_n = div(\vec{g}_n)$ in $D'$, and $\|v_n\|_F = \|g_{1,n}\|_{BMO} + \|g_{2,n}\|_{BMO} \leq C$, we obtain that $\vec{g}_{i,n} \to \vec{g}_i$ in $BMO -$ weak*. We also have, for all $\phi \in D$, $\int_\Omega v_n \phi = -\int_\Omega \vec{g}_{i,n} \cdot \nabla \phi \to -\int_\Omega \vec{g}_i \cdot \nabla \phi$. Therefore $v = div(\vec{g})$ in $D'$.

Since $\|v_n\|_{L^2} \leq C$, up to a subsequence, $v_n \to v$ weakly in $L^2$, and we have $v = div(\vec{g})$ a.e. As $v_n = f - u_n$ and $u_n$ converges to $u$ weakly in $L^2$, we also obtain $v = f - u$ a.e.

By weak* lower semi-continuity, it follows that

$$\|v\|_F \leq \|g_1\|_{BMO} + \|g_2\|_{BMO} \leq \|g_{1,n}\|_{BMO} + \|g_{2,n}\|_{BMO} = \|v_n\|_F,$$
\[ |u|_{BV} \leq \lim \inf |u_n|_{BV}, \]
\[ ||v||_{F} \leq \lim \inf ||v_n||_{F}. \]
Therefore, \( E(u, v) \leq \lim \inf_{n \to \infty} E(u_n, v_n), \) and we obtain existence of minimizers.

5 Approximating the \((BV, F')\) Decomposition Model

Here we do not solve (16) directly, but we adapt the model \([45]\) by adding a fidelity term into the energy. In this decomposition, \( f \in L^2 \) is decomposed into \( u + v + w \), with \( u \in BV, v \in F \), and a small residual \( w \in L^2 \).

The variational problem can be written as

\[
\inf \left\{ E_{\mu}(u, v) = |u|_{BV} + \mu \int_{\Omega} |f - u - v|^2 + \lambda ||v||_{F} : u \in BV, v \in F, \int_{\Omega} u = \int_{\Omega} f \right\}.
\]

Taking \( v = \text{div} (\vec{g}) \), we obtain an equivalent formulation in terms of \( u, g_1, \text{ and } g_2 \):

\[
\inf \left\{ E_{\mu}(u, g_1, g_2) = |u|_{BV} + \mu \int_{\Omega} |f - u - \text{div}(\vec{g})|^2 \\
+ \lambda [||g_1||_{BMO} + ||g_2||_{BMO}] \right\},
\]

where the infimum is taken over \( g_i \in BMO \) and \( u \in BV \) with

\[
\int_{\Omega} u = \int_{\Omega} f.
\]

The existence and uniqueness of a minimizer can be shown for the new model.

**Theorem 2.** Let \( f \in L^2 \). Then there exist \( u \in BV, \) and \( v \in F \cap L^2, \) such that \((u, v)\) solves (19) or (20). If in addition \( \int_{\Omega} f \neq 0 \), then the minimizer is unique.

**Proof:** Existence of minimizers: let \((u_n, v_n)\) be a minimizing sequence of (19) or (20). We have,

\[
|u_n|_{BV} \leq C; \quad ||f - u_n - v_n||_{L^2} \leq C; \quad ||v_n||_{F} \leq C.
\]
From the Poincaré inequality,

$$||u_n - \int_{\Omega} u_n||_{L^2} \leq C||u_n||_{BV}. $$

Since $\int_{\Omega} u_n = \int_{\Omega} f$, $\forall n$, $u_n$ is uniformly bounded in $L^2$. Since $\Omega$ is bounded, $u_n$ is also uniformly bounded in $L^1$. Therefore,

$$||u_n||_{BV} \leq C. \quad (25)$$

Then there exists $u \in BV$, such that, up to a subsequence, $u_n$ converges to $u$ weak* in $BV$. By (23) and uniform boundedness of $u_n$ in $L^2$, $v_n$ is also uniformly bounded in $L^2$. Therefore, there exists $v \in L^2$ such that, up to a subsequence, $v_n$ converges to $v$ weakly in $L^2$.

As $||v_n||_{F} \leq C$, there exists $\bar{g}_n = (g_{1,n}, g_{2,n}) \in BMO(\Omega, \mathbb{R}^2)$, such that $v_n = div(\bar{g}_n)$ in $\mathcal{D}'$, and $||g_{i,n}||_{BMO} \leq C$. Therefore, there exists $g_i \in BMO$, such that $g_{i,n}$ converges to $g_i$ weak* in $BMO$, for $i = 1, 2$. Let $\bar{g} = (g_1, g_2)$.

To show $v = div(\bar{g}) \in F$: let $\varphi \in \mathcal{D}$,

$$\int_{\Omega} v_n \varphi = \int_{\Omega} div(\bar{g}_n) \varphi = -\int_{\Omega} \bar{g}_n \cdot \nabla \varphi. $$

Taking $n \to \infty$, (using weak $L^2$ topology, and weak* $BMO(\Omega, \mathbb{R}^2)$ topology) we obtain

$$\int_{\Omega} v \varphi = -\int_{\Omega} \bar{g} \cdot \nabla \varphi = \int_{\Omega} div(\bar{g}) \varphi. $$

This implies $v = div(\bar{g})$ in $\mathcal{D}'$, as a distribution. But since $v \in L^2$, $v = div(\bar{g})$ a.e. Therefore, $v \in F \cap L^2$.

By weak and weak* lower semi continuity, it follows that

$$|u|_{BV} \leq \lim\inf |u_n|_{BV},$$

$$||f - u - v||_{L^2} \leq \lim\inf ||f - u_n - v_n||_{L^2}, \text{ and}$$

$$||v||_{F} \leq \lim\inf ||v_n||_{F}. $$

Therefore, $\mathcal{E}_\mu(u, v) \leq \mathcal{E}_\mu(u_n, v_n)$, and $(u, v)$ is a minimizer for (19).

Uniqueness of minimizers: denote by $(\hat{u}, \hat{v})$ a minimizer of the energy. Then $\int_{\Omega} \hat{u} = \int_{\Omega} f$.

The energy to be minimized is strictly convex, as the sum of two convex functions ($|u|_{BV} + ||v||_{F}$) and of a strictly convex function $||f - (u + v)||^2_{L^2}$, except in
the direction \((u, -u)\) (as in [7], [6]). Therefore it suffices to check that if \((\hat{u}, \hat{v})\) is a minimizer, then \((\hat{u} + t\hat{u}, \hat{v} - t\hat{u})\) is not a minimizer for \(t \neq 0\). Since \((\hat{u}, \hat{v})\) is a minimizer, then \(\int_{\Omega} \hat{u} = \int_{\Omega} f\). Therefore, if \((\hat{u} + t\hat{u}, \hat{v} - t\hat{u})\) is a minimizer too, then \(\int_{\Omega} (1 + t)\hat{u} = (1 + t) \int_{\Omega} \hat{u} = \int_{\Omega} f\). This is possible only if \(t = 0\), therefore we conclude the uniqueness.

Again as in [7], [6], we can show that the approximated model (19) approaches Meyer’s model (18) as \(\mu \to \infty\). In other words, we have:

**Theorem 3.** Assume \(f \in L^2\) with \(\int_{\Omega} f \neq 0\), and let us assume that problem (18) has a unique solution \((\hat{u}, \hat{v})\). Let us denote by \((u_\mu, v_\mu)\) the unique solution of (19). Then, as \(\mu \to \infty\), \(u_\mu + v_\mu \to f\) strongly in \(L^2\), and \((u_\mu, v_\mu)\) converges to some \((u_0, v_0)\), up to a subsequence. Moreover, \((u_0, v_0) = (\hat{u}, \hat{v})\) is the solution of (18).

**Proof:** First we need to show that there is \(u \in BV\) and \(v \in F\) such that \(\mathcal{E}_\mu(u, v) \leq C\), where \(C\) does not depend on \(\mu\).

From Remark 5, the R-O-F model with some appropriate \(\lambda'\) ensures such a \(u \in BV\) and a \(v \in G \subset F\) such that \(f = u + v\), and \(\int_{\Omega} u = \int_{\Omega} f\). Therefore,

\[
\mathcal{E}_\mu(u, v) = |u|_{BV} + \lambda||v||_F \leq C,
\]

and \(C\) does not depend on \(\mu\). Another pair that satisfies this property is in fact provided by \((\hat{u}, \hat{v})\).

Then, we obtain that

\[
\mathcal{E}_\mu(u_\mu, v_\mu) \leq \mathcal{E}_\mu(u, v) \leq C,
\]

\[
\Rightarrow \mu||f - u_\mu - v_\mu||_{L^2} \leq C,
\]

\[
\Rightarrow ||f - u_\mu - v_\mu||_{L^2} \leq \frac{C}{\mu},
\]

therefore,

\[
||f - u_\mu - v_\mu||_{L^2} \to 0, \text{ as } \mu \to \infty.
\]

Now, as before, we can deduce that \(||u_\mu||_{BV} \leq C\) and \(||v_\mu||_F \leq C\). Then again similarly, we deduce that there is \((u_0, v_0)\) such that, up to a subsequence, \(u_\mu \to u_0\) in \(BV\)-weak* and weakly in \(L^2\), and \(v_\mu \to v_0\) weakly in \(L^2\). Moreover, and as before, we will have that

\[
\mathcal{E}(u_0, v_0) = |u_0|_{BV} + \lambda||v_0||_F \leq |u_\mu|_{BV} + \lambda||v_\mu||_F \leq \mu||f - (u_\mu + v_\mu)||_{L^2}^2 + |u_\mu|_{BV} + \lambda||v_\mu||_F \leq \mathcal{E}_\mu(\hat{u}, \hat{v}) = \mathcal{E}(\hat{u}, \hat{v}),
\]

i.e. \((u_0, v_0) = (\hat{u}, \hat{v})\) the minimizer of (18), and \((u_0, v_0)\) is the limit of \((u_\mu, v_\mu)\) (up to a subsequence), with \(u_0 + v_0 = f\) a.e. in \(\Omega\).
5.1 Characterization of Minimizers

Here we would like to show some properties of minimizers of problem (19), as a generalization of Theorem 3, page 32, [29].

We recall the variational problem of decomposing $f$, via

$$\inf \left\{ \mathcal{E}_\mu(u, v) = |u|_{BV} + \mu \int_\Omega |f - u - v|^2 + \lambda ||v||_F : u \in BV, v \in F \right\}$$ \hspace{1cm} (26)

(note that here we consider a larger space of possible minimizers $(u, v)$, because we do not impose that the mean value of $u$ is equal with the mean value of $f$ a priori; another way would have been to work with the corresponding quotient spaces).

Definition 11. Given a function $w \in L^2(\Omega)$ and $\lambda > 0$, define

$$||w||_{*,\lambda} = \sup_{g \in BV(\Omega), h \in F(\Omega) \cap L^2(\Omega)} \frac{(w, g + h)}{|g|_{BV} + \lambda ||h||_F}, \quad |g|_{BV} + \lambda ||h||_F \neq 0,$$ \hspace{1cm} (27)

where $(\cdot, \cdot)$ is the $L^2$ inner product.

Remark 6. If $\int_\Omega w \neq 0$, then $||w||_{*,\lambda} = \infty$; indeed, we can replace $g$ by $g + c$, with $c \in \mathbb{R}$, and then the supremum will no longer be finite as $|c| \to \infty$.

We have the following characterizations of an optimal decomposition of $f$ using (26), which will be called the $(BV, F)$ model.

Theorem 4. Let $(u, v)$ be an optimal $(BV, F)$ decomposition of $f$, and denote $w = f - u - v$. Then

1. $||f||_{*,\lambda} \leq \frac{1}{2\mu} \Leftrightarrow u = 0, v = 0, \text{ and } w = f.$

2. Suppose $||f||_{*,\lambda} > \frac{1}{2\mu}$, then $(u, v)$ is characterized by the following two conditions

$$||w||_{*,\lambda} = \frac{1}{2\mu}, \text{ and } (w, u + v) = \frac{1}{2\mu}(|u|_{BV} + \lambda ||v||_F).$$ \hspace{1cm} (28)

Proof: The $(BV, F)$ model (26) yields $u = 0$ and $v = 0$ if and only if for any $g \in BV(\Omega)$, $h \in F(\Omega) \cap L^2(\Omega)$,

$$\mu ||f||_{L^2}^2 \leq |g|_{BV} + \mu ||f - g - h||_{L^2}^2 + \lambda ||h||_F.$$ \hspace{1cm} (29)
Equation (29) holds if and only if (by substituting in (29) \( g \) by \( \epsilon g \) and \( h \) by \( \epsilon h \), and taking \( \epsilon \to 0 \)),
\[
|(f, g + h)| \leq \frac{1}{2\mu}(|g|_{BV} + |h||F). \tag{30}
\]
By the definition of \( \|\cdot\|_{*,\lambda} \), we have \( \|f\|_{*,\lambda} \leq \frac{1}{2\mu} \).

For the converse property in 1., assume that \( \|f\|_{*,\lambda} \leq \frac{1}{2\mu} \). Then, for any \( g \in BV(\Omega) \) and \( h \in F(\Omega) \cap L^2(\Omega) \), with \( |g|_{BV} + \lambda |h||F \neq 0 \), we have
\[
(f, g + h) \leq (|g|_{BV} + \lambda |h||F)\|f\|_{*,\lambda} \leq \frac{1}{2\mu}(|g|_{BV} + \lambda |h||F).
\]
We also have
\[
|g|_{BV} + |f - (g + h)|^2_{L^2} + \lambda |h||F = |g|_{BV} + |f|^2_{L^2} - 2\mu (f, g + h) + \mu |g + h|^2_{L^2} + \lambda |h||F \\
\geq |g|_{BV} + |f|^2_{L^2} - (|g|_{BV} + \lambda |h||F) + \mu |g + h|^2_{L^2} + \lambda |h||F \\
= \mu |f|^2_{L^2} + \mu |g + h|^2_{L^2} \geq \mu |f|^2_{L^2} = \mathcal{E}_\mu(0, 0).
\]

Therefore, \( u = 0 \) and \( v = 0 \) gives the optimal decomposition in this case.

Now suppose \( \|f\|_{*,\lambda} > \frac{1}{2\mu} \). Let \((u, v)\) be an optimal \((BV, F)\) decomposition. We have \( u \neq 0 \) or \( v \neq 0 \). For \( g \in BV(\Omega) \), \( h \in F(\Omega) \cap L^2(\Omega) \), and \( \epsilon \in \mathbb{R} \),
\[
|u + \epsilon g|_{BV} + \mu |w - \epsilon (g + h)|^2_{L^2} + \lambda |v + \epsilon h||F \geq |u|_{BV} + |w|_{L^2} + \lambda |v||F, \tag{31}
\]
\[
\Rightarrow |u|_{BV} + |\epsilon| |g|_{BV} + \mu |w - \epsilon (g + h)|^2_{L^2} + \lambda (|v||F + |\epsilon||h||F) \\
\geq |u|_{BV} + |\epsilon| |w|^2_{L^2} + \lambda |v||F, \\
\Rightarrow |\epsilon| |g|_{BV} + \mu |w - \epsilon (g + h)|^2_{L^2} + |\epsilon| |h||F \geq \mu |w|^2_{L^2}, \\
\Rightarrow |\epsilon| |g|_{BV} + \mu (|w|^2_{L^2} - 2\epsilon (w, g + h) + \epsilon^2 |g + h|^2_{L^2}) + \lambda |\epsilon||h||F \geq \mu |w|^2_{L^2}.
\]
Dividing both sides of the last equation by \( \epsilon > 0 \), we obtain
\[
-2\mu (w, g + h) + \epsilon \mu |g + h|^2_{L^2} + |g|_{BV} + \lambda |h||F \geq 0. \tag{32}
\]
Taking \( \epsilon \to 0 \), we obtain
\[
2\mu (w, g + h) \leq |g|_{BV} + \lambda |h||F, \text{ for all } g \in BV(\Omega), h \in F(\Omega) \cap L^2(\Omega).
\]
Therefore,
\[
\|w\|_{*,\lambda} \leq \frac{1}{2\mu}. \tag{33}
\]
If we take $\epsilon \in (-1, 1)$, and replace $(g, h)$ with $(u, v)$ in equation (31), then (31) implies
\[ 2\mu \epsilon (w, u + v) \leq \epsilon (|u|_{BV} + \lambda ||v||_F) + \epsilon^2 \mu ||u + v||^2_{L^2}. \] (34)
If $\epsilon > 0$, $2\mu (w, u + v) \leq (|u|_{BV} + \lambda ||v||_F)$, and if $\epsilon < 0$, $2\mu (w, u + v) \geq (|u|_{BV} + \lambda ||v||_F)$. Therefore equality holds,
\[ (w, u + v) = \frac{1}{2\mu} (|u|_{BV} + \lambda ||v||_F), \] (35)
and (35) together with (33) implies $||w||_{*,\lambda} = \frac{1}{2\mu}$.
Conversely, if (35) holds for some $(u, v)$ and $||w||_{*,\lambda} = \frac{1}{2\mu}$, then for any $g \in BV(\Omega), h \in F(\Omega) \cap L^2(\Omega),$
\[
|u + \epsilon h|_{BV} + \mu ||w - \epsilon(g + h)||^2_{L^2} + \lambda ||v + \epsilon h||_F
\geq 2\mu (w, u + \epsilon g + v + \epsilon h) + \mu ||w||^2_{L^2} - 2\mu \epsilon (w, g + h) + \mu \epsilon^2 ||g + h||^2_{L^2}
= 2\mu (w, u + v) + \mu ||w||^2_{L^2} + \mu \epsilon^2 ||g + h||^2_{L^2}
= |u|_{BV} + \lambda ||v||_F + \mu ||w||^2_{L^2} + \mu \epsilon^2 ||g + h||^2_{L^2}
\geq |u|_{BV} + \lambda ||v||_F + \mu ||w||^2_{L^2}.
\]
Therefore, $(u, v)$ is an optimal $(BV, F)$ decomposition of $f$.

**Remark 7.** Similar results also hold for the optimal $(BV, G)$ decomposition of $f$, with $G$ replacing $F$ in (26).

**Remark 8.** Note that if a given image $f$ does not have zero mean, i.e. $\int_{\Omega} f \neq 0$, then the first property in Theorem 4 will not hold since $||f||_{*,\lambda} = \infty$. Therefore, if $(u_\mu, v_\mu)$ is a minimizer of (26), then $||w_\mu||_{*,\lambda} \to 0$ as $\mu \to \infty$, as expected. Here $w_\mu = f - u_\mu - v_\mu$. This is in agreement with the result from Theorem 3.

### 5.2 Minimization of (20)

For numerical computations, we use $BMO^\beta$, $BMO_D$, and $BMO$ to represent the functions $g_i, i = 1, 2$. From now on we denote $BMO$ to mean either $BMO^\beta$, $BMO_D$, or $BMO$, and $B$ is either an open set, an $\alpha$ translated dyadic square, or a square in $\Omega$. 

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Equation (20) can be further simplified as

\[
\inf \left\{ \mathcal{E}(u, g_1, g_2) = \int_\Omega |\nabla u| + \mu \int_\Omega |f - u - \partial_x g_1 - \partial_y g_2|^2 \right. \\
+ \lambda \left[ \frac{1}{|B_1|} \int_\Omega |g_1 - g_1, B_1| H(\phi_1) \\
+ \frac{1}{|B_2|} \int_\Omega |g_2 - g_2, B_2| H(\phi_2) \right] \right\},
\]

where \( g_{1,B_1} = \frac{\int_\Omega g_1 H(\phi_1)}{\int_\Omega H(\phi_1)} \), \( \phi_1 \) is the level set of \( B_1 \), and \( B_1 \) maximizes \( \|g_1\|_{BMO} \).

The infimum in (36) is taken over all \( u \in BV \), and \( \bar{g} = (g_1, g_2) \) with \( g_i \in BMO \).

Keeping \( B_1 \) and \( B_2 \) fixed for one iteration, and minimizing \( \mathcal{E}(u, g_1, g_2) \) with respect to its variables, we obtain

\[
-\text{div} \left( \frac{\nabla u}{|\nabla u|} \right) - 2\mu(f - u - \partial_x g_1 - \partial_y g_2) = 0, \tag{37}
\]

\[
2\mu \partial_x (f - u - \partial_x g_1 - \partial_y g_2) \\
+ \lambda H(\phi_1) \left[ \frac{g_1 - g_1, B_1}{|g_1 - g_1, B_1|} - \frac{1}{|B_1|} \int_\Omega \frac{g_1 - g_1, B_1}{|g_1 - g_1, B_1|} H(\phi_1) \right] = 0,
\]

\[
2\mu \partial_y (f - u - \partial_x g_1 - \partial_y g_2) \\
+ \lambda H(\phi_2) \left[ \frac{g_2 - g_2, B_2}{|g_2 - g_2, B_2|} - \frac{1}{|B_2|} \int_\Omega \frac{g_2 - g_2, B_2}{|g_2 - g_2, B_2|} H(\phi_2) \right] = 0,
\]

with the boundary conditions on \( \partial \Omega \),

\[
\frac{\nabla u}{|\nabla u|} \cdot \vec{n} = 0 \\
(f - u - \partial_x g_1 - \partial_y g_2) n_x = 0 \\
(f - u - \partial_x g_1 - \partial_y g_2) n_y = 0
\]

where \( \vec{n} = (n_x, n_y) \) is the exterior unit normal to \( \partial \Omega \). At each iteration, for \( g_1 \) and \( g_2 \) fixed or previously estimated, the unknown sets \( B_1 \) are numerically computed and updated by the methods introduced in the previous sections.

Note that by integrating both sides of (37) over \( \Omega \), the constraint in (21) is automatically satisfied.
6 Approximating the \((BV, G)\) Decomposition Model

For comparison with the \((BV, F)\) model, we approximate the \((BV, G)\) model

\[
\inf_{(u,v) \in BV \times G} \{ \mathcal{E}(u,v) = |u|_{BV} + \lambda|v|_G \},
\]

by the model

\[
\begin{align*}
\inf_{(u,\vec{g}) \in BV \times L^\infty(\Omega,\mathbb{R}^2)} \left\{ \mathcal{E}(u, \vec{g}) = |u|_{BV} + \mu \int_\Omega |f - u - \text{div}(\vec{g})|^2 \\
+ \lambda \int_\Omega \sqrt{g_1(x,y)^2 + g_2(x,y)^2} \delta(x-x_0,y-y_0) \right\}.
\end{align*}
\]

(38)

Here \(v = \text{div} \vec{g}, \delta\) is the Dirac function (an impulse function) in two dimensions concentrated at the origin, and

\[
\sqrt{g_1(x_0,y_0)^2 + g_2(x_0,y_0)^2} = ||\vec{g}||_{L^\infty}.
\]

For numerical computation, we approximate \(\delta\) by a smooth version \(\delta_\epsilon\) such that \(\delta_\epsilon \to \delta\), as \(\epsilon \to 0\).

For \((x_0,y_0)\) fixed, but updated at each iteration, and minimizing \(\mathcal{E}(u, \vec{g})\) with respect to \(u, g_1,\) and \(g_2,\) we obtain the following Euler-Lagrange equations

\[
\begin{align*}
-\text{div} \left( \frac{\nabla u}{|\nabla u|} \right) - 2\mu(f - u - \partial_x g_1 - \partial_y g_2) &= 0, \\
2\mu \partial_x(f - u - \partial_x g_1 - \partial_y g_2) + \lambda \frac{g_1}{\sqrt{g_1^2 + g_2^2}} \delta_\epsilon(x_0,y_0) &= 0, \\
2\mu \partial_y(f - u - \partial_x g_1 - \partial_y g_2) + \lambda \frac{g_2}{\sqrt{g_1^2 + g_2^2}} \delta_\epsilon(x_0,y_0) &= 0,
\end{align*}
\]

with the boundary conditions:

\[
\frac{\nabla u}{|\nabla u|} \cdot \vec{n} = 0,
\]

\[
(f - u - \partial_x g_1 - \partial_y g_2)n_x = 0,
\]

\[
(f - u - \partial_x g_1 - \partial_y g_2)n_y = 0.
\]

For approximating the Dirac delta function, we use [13]

\[
\delta_\epsilon(z) = \frac{\epsilon}{\pi(\epsilon^2 + z^2)}.
\]

(39)
7 A More Isotropic \((BV, F)\) Decomposition

Remark 9. \(v = \text{div}(\vec{g})\) is not isotropic. Also on the rectangular grid, \(\text{div}(\frac{\nabla u}{|\nabla u|})\) is zero at horizontal and vertical edges, and nonzero at other edges numerically. More specifically, the vertical and horizontal oscillations are not so well captured in \(v\) if we represent \(v\) as \(\text{div}(\vec{g})\), but are captured in the \(u\) component (as will be seen in our numerical results).

To overcome these effects, we consider a more isotropic decomposition. As in the O-S-V model [34], we impose \(v = \text{div}(\vec{g}) = \Delta P\), for some scalar function \(P\), to allow stronger smoothing on \(u\). Therefore, the model (36) can be rewritten as,

\[
\inf\left\{ \mathcal{E}(u, P) = \int |\nabla u| + \mu \int |f - u - \Delta P|^2 \right. \\
\quad + \lambda \left[ \frac{1}{|B_1|} \int |P_x - P_{x,B_1}| H_\epsilon(\phi_1) \right. \\
\left. + \frac{1}{|B_2|} \int |P_y - P_{y,B_2}| H_\epsilon(\phi_2) \right] \right\},
\]

where \(H_\epsilon\) is a smooth approximation of the Heaviside function \(H\), and the unknown sets \(B_1\) and \(B_2\) maximize the \(BMO\) norms of \(g_1 = P_x\) and of \(g_2 = P_y\). For fixed \(B_1\) and \(B_2\), but updated after each iteration, minimizing \(\mathcal{E}(u, P)\) in (40) with respect to \(u\) and \(P\), we obtain the Euler-Lagrange equations:

\[- \text{div} \left( \frac{\nabla u}{|\nabla u|} \right) - 2\mu (f - u - \Delta P) = 0,\]

\[- 2\mu \Delta (f - u - \Delta P) - \frac{\lambda}{|B_1|} \left[ \partial_x \left( \frac{P_x - P_{x,B_1}}{|P_x - P_{x,B_1}|} H_\epsilon(\phi_1) \right) \right. \\
\left. + \frac{\lambda}{|B_1|^2} \left( \int \frac{P_x - P_{x,B_1}}{|P_x - P_{x,B_1}|} H_\epsilon(\phi_1) \right) \partial_x H_\epsilon(\phi_1) \right. \\
\left. - \frac{\lambda}{|B_2|} \left[ \partial_y \left( \frac{P_y - P_{y,B_2}}{|P_y - P_{y,B_2}|} H_\epsilon(\phi_2) \right) \right. \\
\left. + \frac{\lambda}{|B_2|^2} \left( \int \frac{P_y - P_{y,B_2}}{|P_y - P_{y,B_2}|} H_\epsilon(\phi_2) \right) \partial_y H_\epsilon(\phi_2) \right] = 0,\]
with the boundary conditions:

\[ \frac{\nabla u}{|\nabla u|} \cdot \vec{n} = 0, \]
\[ (f - u - \Delta P)n_x = 0, \]
\[ (f - u - \Delta P)n_y = 0, \]
\[ (\nabla (f - u - \Delta P)) \cdot \vec{n} = 0. \]

8 Numerical Results

In this section, we present numerical results for image denoising and texture decomposition obtained from the proposed models. We also show comparisons with the Vese-Osher model (4), and the Rudin-Osher-Fatemi model (2).

Let \( f \) be the noisy version of the true image \( \bar{u} \) of size \( M \times N \), and \( u \) be the denoised image. Denote,

\[ RMSE = \sqrt{\frac{\sum_{i=1}^{M,N} (u(i,j) - \bar{u}(i,j))^2}{MN}}. \]

We use \( RMSE \) to quantify how good a denoised image is.

Our numerical results obtained use (39) with \( \epsilon = 0.01 \) to approximate the Dirac delta function. We also normalize the image domain \( \Omega \), so that \( \Omega = [0,1] \times [0,1] \).

In Figure 1, we compute the square that maximizes the \( BMO \) norm (6), and then show the contour of the union of the squares (of the same length) such that the right hand side of (6) is within the given percentage of the \( BMO \) norm. We see that the union of the squares captures the most oscillatory regions in the given image.

Figure 2 and Figure 3 show the contours of open sets that optimize the energies (12) and (13), respectively, using the algorithm described in Section 3.1, and the plots showing the evolution of the energies versus the number of iterations.

Figure 4 and Figure 5 show the contours of squares that optimize the energies (14) and (15), respectively, using the algorithm described in Section 3.2, and the plots showing the evolution of the energies versus the number of iterations.

Figure 6 shows the testing images that we use for our experiments.

Figure 7 shows two image denoisings, using the V-O model (4) and the standard R-O-F model (2). The \( RMSE \) for the V-O model is 0.00767165 in 2000 iterations, and the \( RMSE \) for the R-O-F model is 0.00879536.
Figure 8 shows an image denoising using the \((BV, F)\) decomposition model (36), and the plot showing the evolution of the energy (36) with respect to the number of iterations. We use the dyadic \(BMO\) norm in this case. We also obtain similar result with the method described in Section 3.3. The \(RMSE\) for this decomposition is 0.0076569 in 2000 iterations.

Figure 9 shows an image denoising using the \((BV, G)\) decomposition model (38), and the plot showing the evolution of the energy (38) with respect to the number of iterations. We use \(||||\vec{g}|||_{L^\infty}||\) to compute the energy (38). The \(RMSE\) for this decomposition is 0.0077463 in 10000 iterations. This shows that the \((BV, G)\) decomposition has a slower rate of convergence to the steady state in comparison with the \((BV, F)\) decomposition. Notice that we see more the square in the noise component \(f - u\) in the R-O-F model than we see it by the \((BV, F)\) and \((BV, G)\) models.

In Figure 10-15, we show the decomposition of a given image into cartoon and texture components using \((BV, F)\) and \((BV, G)\) models. We remark that both models give very similar results. For the computation of the \(BMO\) norm, we use the dyadic \(BMO\) with a \(\frac{1}{3}\)-translations in Figure 10 and 12, and in Figure 14 we use the algorithm described in Section 3.2 to obtain the optimal square.

Figure 16-18 shows a decomposition using the standard R-O-F model and the proposed models. We remark that the texture parts are better captured in the oscillatory component \(v\) in Figure 18 using the model (40). Here we use the dyadic \(BMO\) with a \(\frac{1}{3}\)-translations.

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Figure 1: The union of the squares within the given percentage of the $BMO$ norm.
Figure 2: Left: An optimal set which gives $||.||_{BMO^\beta}$. Right: The evolution of the energy (12) versus iterations.

Figure 3: Left: An optimal set which gives $||.||_{BMO^\beta}$. Right: The evolution of the square of the energy (13) versus iterations.
Figure 4: Left: An optimal square which gives $||\cdot||_{BMO}$. Right: The evolution of the energy (14) versus iterations.

Figure 5: Left: An optimal square which gives $||\cdot||_{BMO}$. Right: The evolution of the square of the energy (15) versus iterations.
Figure 6: The data images to be decomposed.
Figure 7: Top: A decomposition using V-O with $p = 2$, $RMSE = 0.00767165$. Bottom: A decomposition using R-O-F. $RMSE = 0.008795368$. 
Figure 8: A decomposition using $(BV, F)$, with $v = div(\bar{g})$, and the plot showing the energy (36) versus iterations. $RMSE = 0.0076569$. 

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Figure 9: A decomposition using $(BV, G)$, with $v = \text{div}(\vec{g})$, and the plot showing the evolution of the energy (38) (using $\|\|\vec{g}\|\|_{L^\infty}$) versus iterations. $RMSE = 0.00775965$.  

Figure 10: A decomposition using $(BV, F)$, with $v = \text{div}(\vec{g})$ with translated dyadic BMO.

Figure 11: A decomposition using $(BV, G)$, with $v = \text{div}(\vec{g})$.
Figure 12: A decomposition using $(BV, F)$, with $v = \text{div}(\vec{g})$ with translated dyadic BMO.

Figure 13: A decomposition using $(BV, G)$, with $v = \text{div}(\vec{g})$. 

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Figure 14: A decomposition using $(BV, F)$, with $v = \text{div}(\vec{g})$.

Figure 15: A decomposition using $(BV, G)$, with $v = \text{div}(\vec{g})$. 

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Figure 16: ROF decomposition.

Figure 17: \((BV, G), \text{ with } v = div(\vec{g})\) decomposition.
Figure 18: A decomposition using $(BV, F)$ model (40), with $v = \Delta P$. 

References


